

QUANTUM ERGODIC RESTRICTION FOR CAUCHY DATA: INTERIOR QUE AND RESTRICTED QUE

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ABSTRACT. We prove a quantum ergodic restriction theorem for the Cauchy data of a sequence of quantum ergodic eigenfunctions on a hypersurface H of a Riemannian manifold (M, g) . The technique of proof is to use a Rellich type identity to relate quantum ergodicity of Cauchy data on H to quantum ergodicity of eigenfunctions on the global manifold M . This has the interesting consequence that if the eigenfunctions are quantum unique ergodic on the global manifold M , then the Cauchy data is automatically quantum unique ergodic on H with respect to operators whose symbols vanish to order one on the glancing set of unit tangential directions to H .

1. INTRODUCTION

This article is concerned with the QER (quantum ergodic restriction) problem for hypersurfaces in compact Riemannian manifolds (M, g) . We consider the eigenvalue problem on M

$$\begin{cases} -\Delta_g \varphi_j = \lambda_j^2 \varphi_j, & \langle \varphi_j, \varphi_k \rangle = \delta_{jk} \\ B\varphi_j = 0 \text{ on } \partial M \end{cases},$$

where $\langle f, g \rangle = \int_M f \bar{g} dV$ (dV is the volume form of the metric) and where B is the boundary operator, e.g. $B\varphi = \varphi|_{\partial M}$ in the Dirichlet case or $B\varphi = \partial_\nu \varphi|_{\partial M}$ in the Neumann case. We also allow $\partial M = \emptyset$. We introduce the Planck constant $h_j = \lambda_j^{-1}$; for notational simplicity we often drop the subscript j . We then denote the eigenfunctions in the orthonormal basis by φ_h and the eigenvalues by h^{-2} , so that the eigenvalue problem takes the semi-classical form,

$$(1.1) \quad \begin{cases} (-h^2 \Delta_g - 1)\varphi_h = 0, \\ B\varphi_h = 0 \text{ on } \partial M \end{cases},$$

where $B = I$ or $B = hD_\nu$ in the Dirichlet or Neumann cases respectively. Also, Δ_g denotes the negative Laplacian, e.g. on flat \mathbb{R}^n it denotes $\sum_{j=1}^n \frac{\partial^2}{\partial x_j^2}$.

Let $H \subset M$ be a smooth hypersurface which does not meet ∂M if $\partial M \neq \emptyset$. The main result of this article (Theorem 1) is that the semiclassical Cauchy data

$$(1.2) \quad CD(\varphi_h) := \{(\varphi_h|_H, hD_\nu \varphi_h|_H)\}$$

of eigenfunctions is *always* quantum ergodic along any hypersurface $H \subset M$ if the eigenfunctions are quantum ergodic on the global manifold M . The proof is a generalization of the boundary case where $H = \partial M$, which was proved in [HaZe] and in [Bu]. Our proof is modeled on that of [Bu], developing ideas of [GL] (see also [CTZ2] for an abstract microlocal approach). This automatic QER property of Cauchy data stands in contrast to the conditional nature of the QER property

for the Dirichlet data alone, which requires an “asymmetry” condition on H with respect to geodesics [TZ1, TZ2, DZ]. Note that in the boundary case $H = \partial M$, the Dirichlet resp. Neumann boundary condition kills one of the two components of the Cauchy data, so that the Cauchy data QER theorem appears the same as the QER theorem for Neumann data (resp. Dirichlet data) alone.

As emphasized below the statement of Theorem 1, the deduction of quantum ergodicity of the Cauchy data from quantum ergodicity of the sequence of ambient eigenfunctions holds for the full original sequence. Hence, if the original sequence is a complete orthonormal basis of eigenfunctions, i.e. if Δ_g is QUE (quantum uniquely ergodic), then the Cauchy data of the full orthonormal basis is also QUE on H for any embedded orientable separating hypersurface H . We refer to this as the QUER property. However, it is not necessarily the case that the Cauchy data is QUER for the full algebra of pseudo-differential operators. In Corollary 1.1 it is proved that QUE on (M, g) implies QUER with respect to the subalgebra of semiclassical pseudodifferential operators on H whose symbols vanish to order 1 along S^*H . The restriction on the symbols arises because the passage from QUE in the ambient manifold to QUER on the hypersurface involves multiplying the symbols by a certain factor which vanishes to order one on S^*H , i.e. the unit directions (co-)tangent to H . Therefore, QUE in the ambient manifold does not imply QUER for all pseudodifferential operators on H , and indeed the test operators damp out the possible modes which concentrate microlocally on H . We nevertheless refer to it as a QUER property because it holds for the entire sequence of eigenfunctions; there is no need to remove a subsequence of density zero for the subalgebra limits.

To state the results precisely, we introduce some notation. We work with the semiclassical calculus of pseudo-differential operators as in the references [Bu, DZ, HaZe, TZ1, Zw]; see also §4 for background. On both H and M we fix (Weyl) quantizations $a \rightarrow a^w$ of semi-classical symbols to semi-classical pseudo-differential operators. When it is necessary to indicate which manifold is involved, we either write $Op_H(a)$ for pseudo-differential operators on H or we use capital letters $A^w(x, hD)$ to indicate operators on M and small letters $a^w(y, hD_y)$ to indicate operators on H .

A sequence of functions u_{h_j} on a manifold M indexed by a sequence of Planck constants is said to be quantum ergodic with limit measure $d\mu$ if

$$\langle a^w(x, h_j D_x) u_j, u_j \rangle \rightarrow \omega(a_0) := \int_{T^*M} a_0 d\mu,$$

for all zeroth order semi-classical pseudo-differential operators, where a_0 is the principal symbol of $a^w(x, h_j D_x)$. In the classical case of compact Riemannian manifolds (M, g) with ergodic geodesic flow of [Sch, Zel, CdV, ZZ], the $\{u_j\}$ were a subsequence of density one of an orthonormal basis of eigenfunctions of Δ_g and $d\mu$ is normalized Liouville measure on S^*M . The functional $a \rightarrow \langle a^w(x, h_j D_x) u_j, u_j \rangle$ is often referred to as a microlocal lift (or a Wigner distribution), and the limit measure or state $\omega(a)$ is often called a quantum limit or a semi-classical defect measure. Thus, we use the term ‘quantum ergodic sequence’ in this article to mean a sequence of functions whose microlocal lifts have a unique weak* limit. For the results of this article, it is not necessary to assume that the geodesic flow is ergodic; we show that that Cauchy data of a quantum ergodic sequence of eigenfunctions in

the ambient space M are quantum ergodic on H , no matter what mechanism leads to quantum ergodicity of the original sequence.

We introduce a hypersurface H , which we assume to be orientable, embedded, and separating in the sense that

$$M \setminus H = M_+ \cup M_-$$

where M_{\pm} are domains with boundary in M . This is not a restrictive assumption since we can arrange that any hypersurface is part of the boundary of a domain.

Given a quantization $a \rightarrow Op_H(a)$ of semi-classical symbols $a \in S_{sc}^0(H)$ of order zero (see Appendix 4) to semi-classical pseudo-differential operators on $L^2(H)$, we define the microlocal lifts of the Neumann data as the linear functionals on $a \in S_{sc}^0(H)$ given by

$$\mu_h^N(a) := \int_{B^*H} a d\Phi_h^N := \langle Op_H(a)hD_\nu\varphi_h|_H, hD_\nu\varphi_h|_H \rangle_{L^2(H)}.$$

We also define the *renormalized microlocal lifts* of the Dirichlet data by

$$\mu_h^D(a) := \int_{B^*H} a d\Phi_h^{RD} := \langle Op_H(a)(1 + h^2\Delta_H)\varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)}.$$

Finally, we define the microlocal lift $d\Phi_h^{CD}$ of the Cauchy data to be the sum

$$(1.3) \quad d\Phi_h^{CD} := d\Phi_h^N + d\Phi_h^{RD}.$$

Here, $h^2\Delta_H$ denotes the negative tangential Laplacian for the induced metric on H , so that the operator $(1 + h^2\Delta_H)$ is characteristic precisely on the glancing set S^*H of H . Intuitively, we have renormalized the Dirichlet data by damping out the whispering gallery components.

The distributions μ_h^N, μ_h^D are asymptotically positive, but are not normalized to have mass one and may tend to infinity. They depend on the choice of quantization, but their possible weak* limits as $h \rightarrow 0$ do not, and the results of the article are valid for any choice of quantization. We refer to §4 or to [Zw] for background on semi-classical microlocal analysis.

Our first result is that the Cauchy data of a sequence of quantum ergodic eigenfunctions restricted to H is automatically QER for semiclassical pseudodifferential operators with symbols vanishing on the glancing set S^*H , i.e. that

$$d\Phi_h^{CD} \rightarrow \omega,$$

where

$$\omega(a) = \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi')(1 - |\xi'|^2)^{1/2} d\sigma$$

is the limit state of Theorem 1. This was proved in a different way in [TZ1] in the case of piecewise smooth Euclidean domains. The assumption $H \cap \partial M = \emptyset$ is for simplicity of exposition and because the case $H = \partial M$ is already known.

Theorem 1. *Suppose $H \subset M$ is a smooth, codimension 1 embedded orientable separating hypersurface and assume $H \cap \partial M = \emptyset$. Assume that $\{\varphi_h\}$ is a quantum ergodic sequence of eigenfunctions (1.1). Then the sequence $\{d\Phi_h^{CD}\}$ (1.3) of microlocal lifts of the Cauchy data of φ_h is quantum ergodic on H in the sense that*

for any $a \in S_{sc}^0(H)$,

$$\begin{aligned} & \langle Op_H(a)hD_\nu\varphi_h|_H, hD_\nu\varphi_h|_H \rangle_{L^2(H)} + \langle Op_H(a)(1+h^2\Delta_H)\varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)} \\ & \rightarrow_{h \rightarrow 0^+} \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi') (1 - |\xi'|^2)^{1/2} d\sigma, \end{aligned}$$

where $a_0(x', \xi')$ is the principal symbol of $Op_H(a)$, $-h^2\Delta_H$ is the induced tangential (semiclassical) Laplacian with principal symbol $|\xi'|^2$, μ is the Liouville measure on S^*M , and $d\sigma$ is the standard symplectic volume form on B^*H .

Remark. We emphasize that the limit along H in Theorem 1 holds for the full sequence $\{\varphi_h\}$. Thus, if the full sequence of eigenfunctions is known to be quantum ergodic, i.e. if the sequence is QUE, then the conclusion of the theorem applies to the full sequence of eigenfunctions.

We also remark that although we do not state it formally, Theorem 1 (and indeed all the results in this paper) apply equally well to quasimodes, that is, to approximate eigenfunctions satisfying

$$\|(-h^2\Delta_g - 1)\varphi_h\|_{L^2} = o(h)\|\varphi_h\|_{L^2}$$

as $h \rightarrow 0^+$.

The proof simply relates the interior and restricted microlocal lifts and reduces the QER property along H to the QE property of the ambient manifold. If we assume that QUE holds in the ambient manifold, we automatically get QUER, which is our first Corollary:

Corollary 1.1. *Suppose that $\{\varphi_h\}$ is QUE on M . Then the distributions $\{d\Phi_h^{CD}\}$ have a unique weak* limit*

$$\omega(a) := \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi') (1 - |\xi'|^2)^{1/2} d\sigma.$$

We note that $d\Phi_h^{CD}$ involves the microlocal lift $d\Phi_h^{RD}$ rather than the microlocal lift of the Dirichlet data. However, in Theorem 2, we see that the analogue of Theorem 1 holds for a *density one subsequence* if we use the further renormalized distributions $d\Phi_h^D + d\Phi_h^{RN}$ where the microlocal lift $d\Phi_h^D \in \mathcal{D}'(B^*H)$ of the Dirichlet data of φ_h is defined by

$$\int_{B^*H} a d\Phi_h^D := \langle Op_H(a)\varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)},$$

and

$$\int_{B^*H} a d\Phi_h^{RN} := \langle (1 + h^2\Delta_H + i0)^{-1}Op_H(a)hD_\nu\varphi_h|_H, hD_\nu\varphi_h|_H \rangle_{L^2(H)}.$$

Theorem 2. *Suppose $H \subset M$ is a smooth, codimension 1 embedded orientable separating hypersurface and assume $H \cap \partial M = \emptyset$. Assume that $\{\varphi_h\}$ is a quantum ergodic sequence. Then, there exists a sub-sequence of density one as $h \rightarrow 0^+$ such that for all $a \in S_{sc}^0(H)$,*

$$\begin{aligned} & \langle (1 + h^2\Delta_H + i0)^{-1}Op_H(a)hD_\nu\varphi_h|_H, hD_\nu\varphi_h|_H \rangle_{L^2(H)} + \langle Op_H(a)\varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)} \\ & \rightarrow_{h \rightarrow 0^+} \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi') (1 - |\xi'|^2)^{-1/2} d\sigma, \end{aligned}$$

where $a_0(x', \xi')$ is the principal symbol of $Op_H(a)$.

The additional step in the proof is a pointwise local Weyl law as in [TZ1] section 8.4 showing that only a sparse set of eigenfunctions could scar on the glancing set S^*H . This is precisely the step which is not allowed in the QUER problem. Therefore, QUER for all $Op_H(a)$ might fail for this rescaled problem; to determine whether it holds for all $Op_H(a)$ we would need a new idea. However, the following is a direct consequence of Theorem 2.

Corollary 1.2. *Suppose that $\{\varphi_h\}$ is QUE on M . Then the distributions $\{d\Phi_h^D + d\Phi_h^{RN}\}$ have a unique weak* limit*

$$\omega(a) := \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi') (1 - |\xi'|^2)^{-1/2} d\sigma$$

with respect to the subclass of symbols which vanish on S^*H .

We prove Theorem 1 by means of a Rellich identity adapted from [GL, Bu]. It is also possible to prove the theorem using the layer potential approach in Step 2 (Proof of (7.4)) in section 7 of [HaZe]. To adapt this proof, one would need to introduce a semi-classical Green's function in place of the Euclidean Green's function, verify that it has the properties of the latter in section 4 of [HaZe], and then go through the proof of Step 2. Despite the authors' fondness for the layer potential approach, this proof is much longer than the infinitesimal Rellich identity approach and we have decided to omit the details.

Acknowledgements. The first version of this article was written at the same time as [TZ1, TZ2] but its completion was post-poned while the authors proved the QER phenomenon for Dirichlet data alone. We were further stimulated to complete the article by a discussion with Peter Sarnak at the Spectral Geometry conference at Dartmouth in July, 2010 in which we debated whether QUE in the ambient domain implies QUER along H . We said 'yes', Sarnak said 'no'; Corollaries 1.1 and 1.2 explain the sense in which both answers are right.

The research of H.C. was partially supported by NSF grant # DMS-0900524; J.T. was partially supported by NSERC grant # OGP0170280 and a William Dawson Fellowship; S.Z. was partially supported by NSF grant # DMS-0904252.

2. RELLICH APPROACH: PROOF OF THEOREM 1

We have assumed H is a separating hypersurface, so that H is the boundary of a smooth open submanifold of M , $H = \partial M_+ \subset M$. There is no loss of generality in this assumption, since we may always use a cutoff to a subset of H . We then use a Rellich type identity to write the integral of a commutator over M_+ as a sum of integrals over the boundary (of course the same argument would apply on $M_- = M \setminus M_+$). The argument is partially motivated by Burq's proof of boundary quantum ergodicity (ie. the case $H = \partial M$).

Let $x = (x_1, \dots, x_{n-1}, x_n) = (x', x_n)$ be Fermi normal coordinates in a small tubular neighbourhood $H(\epsilon)$ of H defined near a point $x_0 \in H$. In these coordinates we can locally write

$$H(\epsilon) := \{(x', x_n) \in U \times \mathbb{R}, |x_n| < \epsilon\}.$$

Here $U \subset \mathbb{R}^{n-1}$ is a coordinate chart containing $x_0 \in H$ and $\epsilon > 0$ is arbitrarily small but for the moment, fixed. We let $\chi \in C_0^\infty(\mathbb{R})$ be a cutoff with $\chi(x) = 0$ for

$|x| \geq 1$ and $\chi(x) = 1$ for $|x| \leq 1/2$. In terms of the normal coordinates,

$$-h^2 \Delta_g = \frac{1}{g(x)} h D_{x_n} g(x) h D_{x_n} + R(x_n, x', h D_{x'})$$

where, R is a second-order h -differential operator along H with coefficients that depend on x_n , and $R(0, x', h D_{x'}) = -h^2 \Delta_H$ is the induced tangential semiclassical Laplacian on H .

Let $A(x, h D_x) \in \Psi_{sc}^0(M)$ be an order zero semiclassical pseudodifferential operator on M (see 4). By Green's formula and (1.1) we get the Rellich identity

$$(2.1) \quad \begin{aligned} & \frac{i}{h} \int_{M_+} ([-h^2 \Delta_g, A(x, h D_x)] \varphi_h(x)) \overline{\varphi_h(x)} dx \\ &= \int_H (h D_\nu A(x', x_n, h D_x) \varphi_h|_H) \overline{\varphi_h}|_H d\sigma_H \\ & \quad + \int_H (A(x', x_n, h D_x) \varphi_h|_H) \overline{h D_\nu \varphi_h}|_H d\sigma_H. \end{aligned}$$

Here, $D_{x_j} = \frac{1}{i} \frac{\partial}{\partial x_j}$, $D_{x'} = (D_{x_1}, \dots, D_{x_{n-1}})$, $D_\nu = \frac{1}{i} \partial_\nu$ where ∂_ν is the interior unit normal to M_+ .

Given $a \in S^{0,0}(T^*H \times (0, h_0])$, we then choose

$$A(x', x_n, h D_x) = \chi\left(\frac{x_n}{\epsilon}\right) h D_{x_n} a^w(x', h D').$$

Since $\chi(0) = 1$ it follows that the second term on the right side of (2.1) is just

$$(2.2) \quad \langle a^w(x', h D') h D_{x_n} \varphi_h|_H, h D_{x_n} \varphi_h|_H \rangle.$$

The first term on right hand side of (2.1) equals

$$(2.3) \quad \begin{aligned} & \int_H h D_n (\chi(x_n/\epsilon) h D_n a^w(x', h D') \varphi_h) \Big|_{x_n=0} \overline{\varphi_h} \Big|_{x_n=0} d\sigma_H \\ &= \int_H \left(\chi(x_n/\epsilon) a^w(x', h D') (h D_n)^2 \varphi_h + \frac{h}{i\epsilon} \chi'(x_n/\epsilon) h D_n a^w(x', h D') \varphi_h \right) \Big|_{x_n=0} \overline{\varphi_h} \Big|_{x_n=0} d\sigma_H \\ &= \int_H (\chi(x_n/\epsilon) a^w(x', h D') (1 - R(x_n, x', h D')) \varphi_h) \Big|_{x_n=0} \overline{\varphi_h} \Big|_{x_n=0} d\sigma_H + \mathcal{O}_\epsilon(h), \end{aligned}$$

since $\chi'(0) = 0$ and $((h D_n)^2 + R + O(h)) \varphi_h = \varphi_h$ in these coordinates.

It follows from (2.1)-(2.3) that

$$(2.4) \quad \langle Op_H(a) h D_\nu \varphi_h|_H, h D_\nu \varphi_h|_H \rangle_{L^2(H)} + \langle Op_H(a) (1 + h^2 \Delta_H) \varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)}$$

$$(2.5) \quad = \left\langle \left(\left\{ \xi_n^2 + R(x_n, x', \xi'), \chi\left(\frac{x_n}{\epsilon}\right) \xi_n a(x', \xi') \right\} \right)^w \varphi_h, \varphi_h \right\rangle_{L^2(M_+)} + \mathcal{O}_\epsilon(h).$$

We now assume that φ_h is a sequence of quantum ergodic eigenfunctions, and take the $h \rightarrow 0^+$ limit on both sides of (2.4). We apply interior quantum ergodicity to the term on the right side of (2.4). We compute

$$(2.6) \quad \left\{ \xi_n^2 + R(x_n, x', \xi'), \chi\left(\frac{x_n}{\epsilon}\right) \xi_n a(x', \xi') \right\} = \frac{2}{\epsilon} \chi'\left(\frac{x_n}{\epsilon}\right) \xi_n^2 a(x', \xi') \\ + \chi\left(\frac{x_n}{\epsilon}\right) R_2(x', x_n, \xi'),$$

where R_2 is a zero order symbol. Let $\chi_2 \in C^\infty$ satisfy $\chi_2(t) = 0$ for $t \leq -1/2$, $\chi_2(t) = 1$ for $t \geq 0$, and $\chi_2'(t) > 0$ for $-1/2 < t < 0$, and let ρ be a boundary defining function for M_+ . Then $\chi_2(\rho/\delta)$ is 1 on M_+ and 0 outside a $\delta/2$ neighbourhood. Now the assumptions that the sequence φ_h is quantum ergodic implies that the matrix element of the second term on the right side of (2.6) is bounded by

$$\begin{aligned} & \left| \langle (\chi(x_n/\epsilon) R_2(x, \xi'))^w \varphi_h, \varphi_h \rangle_{L^2(M_+)} \right| \\ & \leq \| \chi_2(\rho/\delta) \chi(x_n/\epsilon) \varphi_h \|_{L^2(M)} \| \tilde{\chi}_2(\rho/\delta) \tilde{\chi}(x_n/\epsilon) \varphi_h \|_{L^2(M)} \\ & = \mathcal{O}_\delta(\epsilon) + o_{\delta, \epsilon}(1), \end{aligned}$$

where $\tilde{\chi}$ and $\tilde{\chi}_2$ are smooth, compactly supported functions which are one on the support of χ and χ_2 respectively. Here, the last line follows from interior quantum ergodicity of the φ_h since the volume of the supports of $\chi(x_n/\epsilon)$ and $\tilde{\chi}(x_n/\epsilon)$ is comparable to ϵ .

To handle the matrix element of the first term on the right side of (2.6), we note that $\chi'(x_n/\epsilon)|_{M_+} = \tilde{\chi}'(x_n/\epsilon)$ for a smooth function $\tilde{\chi} \in C^\infty(M)$ satisfying $\tilde{\chi} = 1$ in a neighbourhood of $M \setminus M_+$ and zero inside a neighbourhood of H . Then, again by interior quantum ergodicity, we have

$$(2.7) \quad 2 \left\langle \left(\frac{1}{\epsilon} \chi'\left(\frac{x_n}{\epsilon}\right) \xi_n^2 a(x', \xi') \right)^w \varphi_h, \varphi_h \right\rangle_{L^2(M_+)} \\ = 2 \left\langle \left(\frac{1}{\epsilon} \tilde{\chi}'\left(\frac{x_n}{\epsilon}\right) \xi_n^2 a(x', \xi') \right)^w \varphi_h, \varphi_h \right\rangle_{L^2(M)} \\ = \frac{2}{\mu(S^*M)} \int_{S^*M} \frac{1}{\epsilon} \tilde{\chi}'\left(\frac{x_n}{\epsilon}\right) (1 - R(x', x_n, \xi')) a(x', \xi') d\mu + O(\epsilon) + o_\epsilon(1) \\ = \frac{2}{\mu(S^*M)} \int_{S^*M_+} \frac{1}{\epsilon} \chi'\left(\frac{x_n}{\epsilon}\right) (1 - R(x', x_n, \xi')) a(x', \xi') d\mu + O(\epsilon) + o_\epsilon(1),$$

since $\tilde{\chi}'$ and χ' are supported inside M_+ . Combining the above calculations yields

$$(2.8) \quad \langle Op_H(a) h D_\nu \varphi_h|_H, h D_\nu \varphi_h|_H \rangle_{L^2(H)} + \langle Op_H(a) (1 + h^2 \Delta_H) \varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)} \\ = \frac{2}{\mu(S^*M)} \int_{S^*M_+} \frac{1}{\epsilon} \chi'\left(\frac{x_n}{\epsilon}\right) (1 - R(x', x_n, \xi')) a(x', \xi') d\mu + O_\delta(\epsilon) + o_{\delta, \epsilon}(1).$$

Finally, we take the $h \rightarrow 0^+$ -limit in (2.8) followed by the $\epsilon \rightarrow 0^+$ -limit, and finally the $\delta \rightarrow 0^+$ limit. The result is that, since the left-hand side in (2.8) is

independent of ϵ and δ ,

$$\begin{aligned}
& \lim_{h \rightarrow 0^+} \langle Op_H(a)hD_\nu \varphi_h|_H, hD_\nu \varphi_h|_H \rangle_{L^2(H)} + \langle Op_H(a)(1+h^2\Delta_H)\varphi_h|_H, \varphi_h|_H \rangle_{L^2(H)} \\
&= \frac{2}{\mu(S^*M)} \int_{S_H^*M} (1 - R(x', x_n = 0, \xi')) d\tilde{\sigma} \\
(2.9) \quad &= \frac{4}{\mu(S^*M)} \int_{B^*H} (1 - |\xi'|^2)^{1/2} a(x', \xi') d\sigma,
\end{aligned}$$

where $d\tilde{\sigma}$ is the symplectic volume form on S_H^*M , and $d\sigma$ is the symplectic volume form on B^*H . \square

3. PROOF OF THEOREM 2 AND COROLLARY 1.2

The proof follows as in Theorem 1 with a few modifications. For fixed $\epsilon_1 > 0$ we choose the test operator

$$(3.1) \quad A(x', x_n, hD_x) := (I + h^2\Delta_H(x', hD') + i\epsilon_1)^{-1} \chi\left(\frac{x_n}{\epsilon}\right) hD_{x_n} a^w(x', hD')$$

and since $WF_h'(\varphi_h|_H) \subset B^*H$ (see [TZ2] section 11) it suffices to assume that $a \in C_0^\infty(T^*H)$ with

$$\text{supp } a \subset B_{1+\epsilon_1}^*(H).$$

Let $\chi_{\epsilon_1}(x', \xi') \in C_0^\infty(B_{1+\epsilon_1}^*(H) \setminus B_{1-2\epsilon_1}^*(H); [0, 1])$ be a cutoff near the glancing set S^*H with $\chi_{\epsilon_1}(x', \xi') = 1$ when $(x', \xi') \in B_{1+\epsilon_1}^*(H) \setminus B_{1-\epsilon_1}^*(H)$. Then, with $A(x, hD_x)$ in (3.1), the same Rellich commutator argument as in Theorem 1 gives

$$\begin{aligned}
(3.2) \quad & \langle (1 + h^2\Delta_H + i\epsilon_1)^{-1} a^w(x', hD')(1 - \chi_{\epsilon_1})^w hD_\nu \varphi_h|_H, hD_\nu \varphi_h|_H \rangle_{L^2(H)} \\
&+ \left\langle a^w(x', hD')(1 - \chi_{\epsilon_1})^w \left(\frac{1 - |\xi'|^2}{1 - |\xi'|^2 + i\epsilon_1} \right)^w \varphi_h|_H, \varphi_h|_H \right\rangle_{L^2(H)} \\
&\rightarrow \frac{4}{\mu(S^*M)} \int_{B^*H} a_0(x', \xi')(1 - \chi_{\epsilon_1}(x', \xi')) \left(\frac{(1 - |\xi'|^2)^{1/2}}{1 - |\xi'|^2 + i\epsilon_1} \right) d\sigma.
\end{aligned}$$

It remains to determine the contribution of the glancing set S^*H . As in [Bu, DZ, HaZe, TZ1] we use a local Weyl law to do this. Because of the additional normal derivative term the argument is slightly different than in the cited articles and so we give some details. For the rest of this proof, we need to recall that $h \in \{\lambda_j^{-1}\}$, and we write h_j for this sequence to emphasize that it is a discrete sequence of values $h_j \rightarrow 0$. Since $\|a^w(x', hD')\|_{L^2 \rightarrow L^2} = O(1)$, it follows that for $h \in (0, h_0(\epsilon_1)]$ with $h_0(\epsilon_1) > 0$ sufficiently small,

$$\begin{aligned}
(3.3) \quad & \frac{1}{N(h)} \sum_{h_j \geq h} \left| \langle a^w(x', hD') \chi_{\epsilon_1}^w \varphi_{h_j}|_H, \varphi_{h_j}|_H \rangle_{L^2(H)} \right| \\
&\leq C \frac{1}{N(h)} \sum_{h_j \geq h} (\|\chi_{\epsilon_1}^w \varphi_{h_j}|_H\|_{L^2(H)} \|\chi_{2\epsilon_1}^w \varphi_{h_j}|_H\|_{L^2(H)} + \mathcal{O}(h_j^\infty)) \\
&\leq \frac{C}{2} \frac{1}{N(h)} \sum_{h_j \geq h} (\|\chi_{\epsilon_1}^w \varphi_{h_j}|_H\|_{L^2(H)}^2 + \|\chi_{2\epsilon_1}^w \varphi_{h_j}|_H\|_{L^2(H)}^2 + \mathcal{O}(h_j^\infty)) \\
&= \mathcal{O}(\epsilon_1^2).
\end{aligned}$$

By a Fourier Tauberian argument [TZ1] section 8.4, it follows that for $h \in (0, h_0(\epsilon_1)]$

$$(3.4) \quad \frac{1}{N(h)} \sum_{h_j \geq h} |\chi_{\epsilon_1, 2\epsilon_1}^w \varphi_{h_j}|_H(x')|^2 = \mathcal{O}(\epsilon_1^2)$$

uniformly for $x' \in H$. The last estimate in (3.3) follows from (3.4) by integration over H .

To estimate the normal derivative terms, we first recall the standard resolvent estimate

$$\|(1 + h^2 \Delta_H + i\epsilon_1)^{-1} u\|_{H_h^2(H)} \leq C \epsilon_1^{-1} \|u\|_{L^2(H)},$$

where H_h^2 is the semiclassical Sobolev space of order 2 (see [Zw] Lemma 13.6). Applying the obvious embedding $H_h^2(H) \subset L^2(H)$, we recover

$$\begin{aligned} \|(1 + h^2 \Delta_H + i\epsilon_1)^{-1} u\|_{L^2(H)} &\leq C \|(1 + h^2 \Delta_H + i\epsilon_1)^{-1} u\|_{H_h^2(H)} \\ &\leq C \epsilon_1^{-1} \|u\|_{L^2(H)} \end{aligned}$$

to get that

$$(3.5) \quad \begin{aligned} &\frac{1}{N(h)} \sum_{h_j \geq h} | \langle (1 + h^2 \Delta_H + i\epsilon_1)^{-1} a^w(x', hD') \chi_{\epsilon_1}^w h_j D_{x_n} \varphi_{h_j}|_H, h_j D_{x_n} \varphi_{h_j}|_H \rangle_{L^2(H)} | \\ &\leq C' \epsilon_1^{-1} \frac{1}{N(h)} \sum_{h_j \geq h} (\|\chi_{\epsilon_1}^w h D_{x_n} \varphi_{h_j}|_H\|_{L^2(H)} \|\chi_{2\epsilon_1}^w h D_{x_n} \varphi_{h_j}|_H\|_{L^2(H)} + \mathcal{O}(h_j^\infty)) \\ &\leq \frac{C' \epsilon_1^{-1}}{2} \frac{1}{N(h)} \sum_{h_j \geq h} (\|\chi_{\epsilon_1}^w h D_{x_n} \varphi_{h_j}|_H\|_{L^2(H)}^2 + \|\chi_{2\epsilon_1}^w h D_{x_n} \varphi_{h_j}|_H\|_{L^2(H)}^2 + \mathcal{O}(h_j^\infty)) \\ &= \mathcal{O}(\epsilon_1^{-1} \epsilon_1^2) \\ &= \mathcal{O}(\epsilon_1). \end{aligned}$$

The last estimate follows again from the Fourier Tauberian argument in [TZ1] section 8.4, which gives

$$(3.6) \quad \frac{1}{N(h)} \sum_{h_j \geq h} |\chi_{\epsilon_1, 2\epsilon_1}^w h_j D_{x_n} \varphi_{h_j}|_H(x')|^2 = \mathcal{O}(\epsilon_1^2)$$

uniformly for $x' \in H$.

Since $\epsilon_1 > 0$ is arbitrary, Theorem 2 follows from (3.3) and (3.5) by letting $\epsilon_1 \rightarrow 0^+$ in (3.2). □

3.1. Proof of Corollary 1.2. We now observe that Corollary 1.2 follows almost immediately from the proof of Theorem 2. To see this, we notice that by restricting our attention to symbols which vanish on the glancing set, we do not need to pass through the local Weyl law/Tauberian argument, which is the step by which one extracts a density one subsequence. Hence the result applies to the full sequence. □

4. APPENDIX

4.1. Semiclassical symbols. Let M be a compact manifold. By a semiclassical symbol $a \in S^{m,k}(T^*M \times [0, h_0))$, we mean a smooth function possessing an asymptotic expansion as $h \rightarrow 0$ of the form,

$$(4.1) \quad a(x, \xi, h) \sim_{h \rightarrow 0^+} \sum_{j=0}^{\infty} a_{k-j}(x, \xi) h^{m+j},$$

with $a_{k-j} \in S_{1,0}^k(T^*M)$. Here, $S_{1,0}^k$ is the standard Hörmander class consisting of smooth functions $a(x, \xi)$ satisfying the estimates $|\partial_x^\alpha \partial_\xi^\beta a(x, \xi)| \leq C_{\alpha,\beta} \langle \xi \rangle^{k-|\beta|}$ for all multi-indices $\alpha, \beta \in N^n$. We say that $A(h) \in Op_h^w(S^{m,k}(T^*M \times [0, h_0)))$ provided its Schwartz kernel is locally of the form

$$(4.2) \quad A(h)(x, y) = (2\pi h)^{-n} \int_{\mathbb{R}^n} e^{i\langle x-y, \xi \rangle / h} a\left(\frac{x+y}{2}, \xi, h\right) d\xi$$

with $a \in S^{m,k}$. We denote the operator $A(h)$ by $a^w(x, hD_x)$ (or simply a^w). By a symbol of order zero we mean that $a \in S^{0,0}$, and we refer to $a_0(x, \xi)$ as the principal symbol. In the latter case, we simply write

$$S_{sc}^0(M) := S^{0,0}, \quad \Psi_{sc}^0(M) := Op_h^w(S^{0,0}).$$

Finally, when $H \subset M$ is a hypersurface and $a \in S_{sc}^0(H)$, we sometimes write $Op_H(a) = a^w$ to indicate dependence on the submanifold, H . We refer to [Zw] for background.

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