CUTOFF RESOLVENT ESTIMATES AND THE SEMILINEAR SCHRÖDINGER EQUATION

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ABSTRACT. This paper shows how abstract resolvent estimates imply local smoothing for solutions to the Schrödinger equation. If the resolvent estimate has a loss when compared to the optimal, non-trapping estimate, there is a corresponding loss in regularity in the local smoothing estimate. As an application, we apply well-known techniques to obtain well-posedness results for the semi-linear Schrödinger equation.

1. INTRODUCTION

In this short note we show how cutoff semiclassical resolvent estimates for the Laplacian on a non-compact manifold, with spectral parameter on the real axis, lead to well-posedness results for the semilinear Schrödinger equation. Motivated by the requirements of [Chr3] and [BGT2], and the microlocal inverse estimates of [Chr1, Chr2], we first prove a general theorem for a large class of resolvents. Following the recent work of Nonnenmacher-Zworski [NoZw], we apply the general theorem in the case there is a hyperbolic fractal trapped set.

Let (M, g) be a Riemannian manifold of dimension n without boundary, with (non-negative) Laplace-Beltrami operator $-\Delta$ acting on functions. The Laplace-Beltrami operator is an unbounded, essentially self-adjoint operator on $L^2(M)$ with domain $H^2(M)$. We assume (M, g) is asymptotically Euclidean in the sense of [NoZw, (3.7)-(3.9)] and that the classical resolvent $(-\Delta - (\lambda^2 + i\epsilon))^{-1}$ obeys a limiting absorption principle as $\epsilon \to 0+$, $\lambda \neq 0$.

Our first result is that if we have cutoff semiclassical resolvent estimates with a sufficiently small loss, then we have weighted smoothing for the Schrödinger propagator with a loss. Let ρ_s be a smooth, non-vanishing weight function satisfying

(1.1)
$$\rho_s(x) \equiv \left\langle d_g(x, x_0) \right\rangle^{-s},$$

for some fixed x_0 and x outside a compact set.

Theorem 1. Suppose for each compactly supported function $\chi \in C_c^{\infty}(M)$ with sufficiently small support, there is $h_0 > 0$ such that the semi-classical Laplace-Beltrami operator satisfies

(1.2)
$$\|\chi(-h^2\Delta - E)^{-1}\chi u\|_{L^2(M)} \le \frac{g(h)}{h} \|u\|_{L^2(M)}, \ E > 0$$

uniformly in $0 < h \leq h_0$, where $g(h) \geq c_0 > 0$, $g(h) = o(h^{-1})$. Then for each T > 0 and s > 1/2, there is a constant $C = C_{T,s} > 0$ such that

(1.3)
$$\int_0^T \left\| \rho_s e^{it\Delta} u_0 \right\|_{H^{1/2-\eta}(M)}^2 dt \le C \| u_0 \|_{L^2(M)}^2,$$

where $\eta \geq 0$ satisfies

(1.4)
$$g(h)h^{2\eta} = \mathcal{O}(1).$$

and ρ_s is given by (1.1).

The assumption that (M, g) is asymptotically Euclidean is that there exists $R_0 > 0$ sufficiently large that, on each infinite branch of $M \setminus B(0, R_0)$, the semiclassical Laplacian $-h^2\Delta$ takes the form

$$-h^2\Delta|_{M\setminus B(0,R_0)} = \sum_{|\alpha| \le 2} a_{\alpha}(x,h)(hD_x)^{\alpha},$$

with $a_{\alpha}(x,h)$ independent of h for $|\alpha| = 2$,

$$\sum_{\substack{|\alpha|=2\\ |\alpha|\leq 2}} a_{\alpha}(x,h)(hD_x)^{\alpha} \ge C^{-1}|\xi|^2, \ 0 < C < \infty, \text{ and}$$
$$\sum_{\substack{|\alpha|\leq 2}} a_{\alpha}(x,h)(hD_x)^{\alpha} \to |\xi|^2, \text{ as } |x| \to \infty \text{ uniformly in } h.$$

In order to quote the results of [NoZw] we also need the following analyticity assumption: $\exists \theta_0 \in [0, \pi)$ such that the $a_{\alpha}(x, h)$ are extend holomorphically to

 $\{r\omega: \omega \in \mathbb{C}^n, \text{ dist} (\omega, \mathbb{S}^n) < \epsilon, \ r \in \mathbb{C}, \ |r| \ge R_0, \ \arg r \in [-\epsilon, \theta_0 + \epsilon)\}.$

As in [NoZw], the analyticity assumption immediately implies

$$\partial_x^\beta \left(\sum_{|\alpha| \le 2} a_\alpha(x,h) \xi^\alpha - |\xi|^2 \right) = o(|x|^{-|\beta|}) \left\langle \xi \right\rangle^2, \ |x| \to \infty.$$

Recall the free Laplacian $(-\Delta_0 - \lambda^2)^{-1}$ on \mathbb{R}^n has a holomorphic continuation from $\operatorname{Im} \lambda > 0$ to $\lambda \in \mathbb{C}$ for $n \geq 3$ odd, and to the logarithmic covering space for neven. This motivates the limiting absorption assumption, that

$$\lim_{\epsilon \to 0+, \ \lambda \neq 0} \rho_s (-\Delta - (\lambda^2 + i\epsilon))^{-1} \rho_s$$

exists as a bounded operator

$$L^2(M, d\mathrm{vol}_g) \to L^2(M, d\mathrm{vol}_g),$$

provided s > 1/2. As in the free case, we allow a possible logarithmic singularity at $\lambda = 0$.

The problem of "local smoothing" estimates for the Schrödinger equation has a long history. The sharpest results to date are those of Doi [Doi] and Burq [Bur]. Doi proved if M is asymptotically Euclidean, then one has the estimate

(1.5)
$$\int_0^T \left\| \chi e^{it\Delta} u_0 \right\|_{H^{1/2}(M)}^2 dt \le C \| u_0 \|_{L^2(M)}^2$$

for $\chi \in \mathcal{C}^{\infty}_{c}(M)$ if and only if there are no trapped sets. Burq's paper showed if there is trapping due to the presence of several convex obstacles in \mathbb{R}^{n} satisfying certain assumptions, then one has the estimate (1.5) with the $H^{1/2}$ norm replaced by $H^{1/2-\eta}$ for $\eta > 0$. In [Chr3], the author considered an arbitrary, single trapped hyperbolic orbit. One of the goals of this paper is to use estimates obtained by Nonnenmacher-Zworski [NoZw] for fractal hyperbolic trapped sets to obtain similar results to [Chr3] for the semilinear Schrödinger equation. To that end we have the following corollary to Theorem 1. **Corollary 1.1.** Assume (M, g) admits a hyperbolic fractal trapped set, K_E , in the energy level E > 0 and that the topological pressure $P_E(1/2) < 0$. Then $-h^2\Delta - E$ satisfies (1.2) for some E > 0 with $g(h) = C \log(1/h)$, and for every $\eta > 0$, T > 0, and s > 1/2, there exists a constant $C = C_{P_E,\eta,T,s} > 0$ such that

$$\int_0^T \left\| \rho_s e^{it\Delta} u_0 \right\|_{H^{1/2-\eta}(M)}^2 dt \le C \| u_0 \|_{L^2(M)}^2.$$

We remark that the assumption $P_E(1/2) < 0$ implies the trapped set K_E is filamentary or "thin" (see [NoZw] for definitions).

We consider the following semilinear Schrödinger equation problem:

(1.6)
$$\begin{cases} i\partial_t u + \Delta u = F(u) \text{ on } I \times M; \\ u(0,x) = u_0(x), \end{cases}$$

where $I \subset \mathbb{R}$ is an interval containing 0. Here the nonlinearity F satisfies

$$F(u) = G'(|u|^2)u,$$

and $G: \mathbb{R} \to \mathbb{R}$ is at least C^3 and satisfies

$$|G^{(k)}(r)| \le C_k \langle r \rangle^{\beta-k},$$

for some $\beta \geq \frac{1}{2}$.

In §3 we prove a family of Strichartz-type estimates which will result in the following well-posedness theorem.

Theorem 2. Suppose (M, g) satisfies the assumptions of the introduction, and set

(1.7)
$$\delta = \frac{4\eta}{2\eta + 1} \ge 0$$

Then for each

(1.8)
$$s > \frac{n}{2} - \frac{2}{\max\{2\beta - 2, 2\}} + \delta$$

and each $u_0 \in H^s(M)$ there exists $p > \max\{2\beta - 2, 2\}$ and $0 < T \le 1$ such that (1.6) has a unique solution

(1.9)
$$u \in C([-T,T]; H^s(M)) \cap L^p([-T,T]; L^{\infty}(M)).$$

Moreover, the map $u_0(x) \mapsto u(t,x) \in C([-T,T]; H^s(M))$ is Lipschitz continuous on bounded sets of $H^s(M)$, and if $||u_0||_{H^s}$ is bounded, T is bounded from below.

If, in addition, (M,g) satisfies the assumptions of Corollary 1.1, $n \leq 3$, $\beta < 3$, and $G(r) \rightarrow +\infty$ as $r \rightarrow +\infty$, then u in (1.9) extends to a solution

$$u \in C((-\infty,\infty); H^1(M)) \cap L^p((-\infty,\infty); L^\infty(M)).$$

Remark 1.2. In particular, the cubic defocusing non-linear Schrödinger equation is globally H^1 -well-posed in three dimensions with a fractal trapped hyperbolic set which is sufficiently filamentary. Of course other nonlinearities can be considered, but for simplicity we consider only these in this work.

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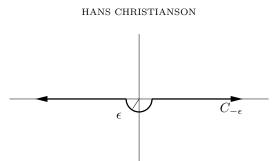


FIGURE 1. The curve $C_{-\epsilon}$ in the complex plane.

2. Proof of Theorem 1

Since we are assuming $(-\Delta - z)^{-1}$ obeys a limiting absorption principle, we have $\|\rho_s(-\Delta - (\tau - i\epsilon))^{-1}\rho_s\|_{L^2 \to L^2} \leq C_\epsilon$

for $0 < \epsilon_0 \le |\tau| \le C$. For $|\sigma| \ge C$ for some C > 0, $\sigma \in \mathbb{C}$ in a neighbourhood of the real axis, write

$$\begin{aligned} -\Delta - \sigma &= -\Delta - \frac{z}{h^2} \\ &= h^{-2} (-h^2 \Delta - z), \end{aligned}$$

for

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$$z \in [E - \alpha, E + \alpha] + i[-c_0h, c_0h].$$

Now

$$(-h^2\Delta - z)$$

is a Fredholm operator for z in the specified range, and hence the "gluing" techniques from [Vod] and [Chr3, §2] can be used to conclude for s > 1/2,

$$\rho_s(-h^2\Delta-z)^{-1}\rho_s$$

has a holomorphic extension to a slightly smaller neighbourhood in $\boldsymbol{z},$ and in particular,

$$\|\rho_s(-h^2\Delta - E)^{-1}\rho_s\|_{L^2 \to L^2} \le C \frac{g(h)}{h}.$$

Rescaling, we have

(2.1)
$$\left\|\rho_s(-\Delta-\tau)^{-1}\rho_s\right\|_{L^2\to L^2} \le C \frac{g(\langle\tau\rangle^{1/2})}{\langle\tau\rangle^{1/2}}, \ \tau\in\mathcal{C}_{\pm\epsilon},$$

where (see Figure 1)

$$\mathcal{C}_{\pm\epsilon} = \{\tau \in \mathbb{R} : |\tau| \ge \epsilon\} \cup \{\tau \in \mathbb{C} : |\tau| = \epsilon, \pm \operatorname{Im} \tau \ge 0\}$$

As in [Chr3] and [Bur], the following lemma follows from integration by parts and interpolation, together with the condition on η , (1.4).

Lemma 2.1. With the notation and assumptions above, we have

$$\|\rho_s(-\Delta-\tau)^{-1}\rho_s\|_{L^2\to H^1} \le Cg(\langle\tau\rangle^{1/2}), \ \tau \in \mathcal{C}_{\pm\epsilon},$$

and for every $r \in [-1, 1]$,

$$\|\rho_s(-\Delta-\tau)^{-1}\rho_s\|_{H^r\to H^{1+r-\eta/2}} \le C, \ \tau \in \mathcal{C}_{\pm\epsilon}.$$

Theorem 1 now follows from the standard " TT^* " argument, letting $\epsilon \to 0$ in (2.1) (see [BGT2], the references cited therein, and [Chr3]).

The following Corollary uses interpolation with an H^2 estimate to replace the $H^{1/2-\eta}$ norm on the left hand side of (1.3) with $H^{1/2}$, and will be of use in §3. See [Chr3] for the details of the proof.

Corollary 2.2. Suppose (M, g) satisfies the assumptions of Theorem 1. For each T > 0 and s > 1/2, there is a constant C > 0 such that

(2.2)
$$\int_0^T \left\| \rho_s e^{it\Delta} u_0 \right\|_{H^{1/2}(M)}^2 dt \le C \| u_0 \|_{H^{\delta}(M)}^2,$$

where $\delta \geq 0$ is given by (1.7).

In particular, if (M,g) satisfies the assumptions of Corollary 1.1, then for any $\delta > 0$, there is $C = C_{\delta} > 0$ such that (2.2) holds.

3. Strichartz-type Inequalities

In this section we give several families of Strichartz-type inequalities and prove Theorem 2. The statements and proofs are mostly adaptations of similar inequalities in [BGT2], so we leave out the proofs of these in the interest of space.

If we view $M \setminus U$, where U is a neighbourhood of K_E , as a manifold with nontrapping geometry, we may apply the results of [HTW] or [BoTz] to a solution of the Schrödinger equation away from the trapping region, resulting in perfect Strichartz estimates. For this section we need (1.3) only with a compact cutoff χ instead of with the more general weight ρ_s .

Proposition 3.1. For every $0 < T \leq 1$ and each $\chi \in C_c^{\infty}(M)$ satisfying $\chi \equiv 1$ near U, there is a constant C > 0 such that

(3.1)
$$\|(1-\chi)u\|_{L^p([0,T])W^{s,q}(M)} \le C \|u_0\|_{H^s(M)},$$

where $u = e^{it\Delta}u_0$, $s \in [0, 1]$, and (p, q), p > 2 satisfy

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}.$$

Remark 3.2. In the sequel, wherever unambiguous, we will write

$$L^{p}_{T}W^{s,q} := L^{p}([0,T])W^{s,q}(M)$$

and

$$H^s := H^s(M).$$

Proposition 3.3. Suppose (M,g) satisfies the assumptions of the Introduction, $u = e^{it\Delta}u_0$, and

$$v = \int_0^t e^{i(t-\tau)\Delta} f(\tau) d\tau.$$

Then for each $0 < T \le 1$ and $\delta \ge 0$ satisfying (1.7), we have the estimates (3.2) $\|u\|_{L^p_m W^{s-\delta,q}} \le C \|u_0\|_{H^s}$

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and

(3.3)
$$\|v\|_{L^p_T W^{s-\delta,q}} \le C \|f\|_{L^1_T H^s},$$

where $s \in [0, 1]$ and (p, q), p > 2 satisfy the Euclidean scaling

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}$$

The proof uses a local WKB expansion localized also in time to the scale of inverse frequency, followed by summing over frequency bands (see [Chr3] and [BGT1]). The only difference here is the explicit dependence of δ on η , which is related to the growth of the function g(h).

Proof of Theorem 2. The proof of Theorem 2 is a slight modification of the proof of Proposition 3.1 in [BGT1], but we include it here in the interest of completeness. Fix s satisfying 1.8 and choose $p > \max\{2\beta - 2, 2\}$ satisfying

$$s > \frac{n}{2} - \frac{2}{p} + \delta \ge \frac{n}{2} - \frac{1}{\max\{2\beta - 2, 2\}}$$

where $\delta \geq 0$ satisfies (1.7). Set $\sigma = s - \delta$ and

$$Y_T = C([-T, T]; H^s(M)) \cap L^p([-T, T]; W^{\sigma, q}(M))$$

for

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2},$$

equipped with the norm

$$\|u\|_{Y_T} = \max_{|t| \le T} \|u(t)\|_{H^s(M)} + \|u\|_{L^p_T W^{\sigma,q}}.$$

Let Φ be the nonlinear functional

$$\Phi(u) = e^{it\Delta}u_0 - i\int_0^t e^{i(t-\tau)\Delta}F(u(\tau))d\tau.$$

If we can show that $\Phi: Y_T \to Y_T$ and is a contraction on a ball in Y_T centered at 0 for sufficiently small T > 0, this will prove the first assertion of the Proposition, along with the Sobolev embedding

(3.5)
$$W^{\sigma,q}(M) \subset L^{\infty}(M),$$

since $\sigma > n/q$. From Proposition 3.3, we bound the W^{σ} part of the Y_T norm by the H^s norm, giving

$$\begin{split} \|\Phi(u)\|_{Y_{T}} &\leq C\left(\|u_{0}\|_{H^{s}} + \int_{-T}^{T} \|F(u(\tau))\|_{H^{s}} d\tau\right) \\ &\leq C\left(\|u_{0}\|_{H^{s}} + \int_{-T}^{T} \|(1 + |u(\tau)|)\|_{L^{\infty}}^{2\beta - 2})\|u(\tau)\|_{H^{s}} d\tau\right), \end{split}$$

where the last inequality follows by our assumptions on the structure of F. Applying Hölder's inequality in time with $\tilde{p} = p/(2\beta - 2)$ and \tilde{q} satisfying

$$\frac{1}{\tilde{q}} + \frac{1}{\tilde{p}} = 1$$

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gives

$$\|\Phi(u)\|_{Y_T} \le C\left(\|u_0\|_{H^s} + T^{\gamma}\|u\|_{L^{\infty}_T H^s}\|(1+|u|)\|_{L^p_T L^{\infty}}^{2\beta-2}\right)$$

where $\gamma = 1/\tilde{q} > 0$. Thus

$$\|\Phi(u)\|_{Y_T} \le C\left(\|u_0\|_{H^s} + T^{\gamma}(\|u\|_{Y_T} + \|u\|_{Y_T}^{2\beta})\right).$$

Similarly, we have for $u, v \in Y_T$,

$$(3.6) \|\Phi(u) - \Phi(v)\|_{Y_T} \le (3.7) \le CT^{\gamma} \|u - v\|_{L^{\infty}_T H^s} \|(1 + |u|)\|_{L^p_T L^{\infty}}^{2\beta - 2} + \|(1 + |v|)\|_{L^p_T L^{\infty}}^{2\beta - 2})$$

$$\leq CT^{\gamma} \|u - v\|_{Y_T} \|(1 + |u|)\|_{Y_T}^{2\beta - 2} + \|(1 + |v|)\|_{Y_T}^{2\beta - 2}),$$

which is a contraction for sufficiently small T. This concludes the proof of the first assertion in the Proposition.

To get the second assertion, we observe from 3.6 and the definition of Y_T , if u and v are two solutions to (1.6) with initial data u_0 and u_1 respectively, so

$$\widetilde{\Phi}(v) = e^{it\Delta}u_1 - i\int_0^t e^{i(t-\tau)\Delta}F(v(\tau))d\tau,$$

we have

$$\begin{aligned} \max_{\substack{|t| \leq T}} \|u(t) - v(t)\|_{H^s} \\ &= \max_{\substack{|t| \leq T}} \|\Phi(u)(t) - \widetilde{\Phi}(v)(t)\|_{H^s} \\ &\leq C \quad \left(\|u_0 - u_1\|_{H^s} \\ &+ T^{\gamma} \max_{\substack{|t| \leq T}} \|u(t) - v(t)\|_{H^s} \|(1 + |u|)\|_{L^p_T L^{\infty}}^{2\beta - 2} + \|(1 + |v|)\|_{L^p_T L^{\infty}}^{2\beta - 2} \right), \end{aligned}$$

which, for T > 0 sufficiently small gives the Lipschitz continuity.

If (M, g) satisfies the assumptions of Corollary 1.1, $n \leq 3$, $\beta < 3$, and $G(r) \rightarrow +\infty$ as $r \rightarrow +\infty$, we can take s and p satisfying $p > \max\{2\beta - 2, 2\}$ and

$$s > \frac{n}{2} - \frac{2}{p} + \delta \ge \frac{n}{2} - \frac{2}{\max\{2\beta - 2, 2\}}$$

for any $\delta > 0$. Then $\sigma = s - \delta > q/n$ and the preceding argument holds. Finally, the proof of the global well-posedness now follows from the standard global well-posedness arguments from, for example, [Caz, Chapter 6].

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